

Leptonic decay-constant ratio f_K/f_π from clover-improved $N_f = 2 + 1$ QCD

Enno E. Scholz*

Universität Regensburg, Institut für Theoretische Physik, D-93040 Regensburg, Germany
E-mail: enno.scholz@physik.uni-regensburg.de

Stephan Dürr

Bergische Universität Wuppertal, Gaußstr. 20, D-42119 Wuppertal, Germany
Jülich Supercomputing Centre, Forschungszentrum Jülich, D-52425 Jülich, Germany
E-mail: durr (AT) itp.unibe.ch

The leptonic decay-constant ratio f_K/f_π is calculated from lattice-QCD simulations using $N_f = 2 + 1$ dynamical fermion flavors in the clover-improved formulation and 2-HEX smearing. The simulations were performed for a range of mass-degenerate light quarks including the physical point and at various lattice couplings and volumes, allowing to quantify all relevant sources of systematic uncertainties for our final number of the decay-constant ratio. Utilizing input from chiral perturbation theory, we also quote the charged decay-constant ratio f_{K^\pm}/f_{π^\pm} . With further input from super-allowed nuclear β -decays, eventually we obtain an estimate for the CKM-matrix element V_{us} .

34th annual International Symposium on Lattice Field Theory
24-30 July 2016
University of Southampton, UK

*Speaker.

1. Introduction

In order to determine the ratio V_{us}/V_{ud} of Cabibbo-Kobayashi-Maskawa (CKM) matrix elements¹ from experimentally measured decay widths of the leptonic kaon and pion decays according to [1]

$$\frac{\Gamma(K^\pm \rightarrow \ell \nu_\ell)}{\Gamma(\pi^\pm \rightarrow \ell \nu_\ell)} = \frac{V_{us}^2}{V_{ud}^2} \frac{f_{K^\pm}^2}{f_{\pi^\pm}^2} \frac{M_{K^\pm}^2}{M_{\pi^\pm}^2} \frac{(1 - m_\ell^2/M_{K^\pm}^2)^2}{(1 - m_\ell^2/M_{\pi^\pm}^2)^2} (1 + \delta_{\text{em}}) \quad (1.1)$$

one needs to know—besides the radiative correction δ_{em} —the ratio of the decay constants f_{K^\pm}/f_{π^\pm} , i.e. the ratio of the hadronic form factors for these decay channels. Here we will report on a calculation of this decay-constant ratio obtained in the isospin symmetric limit of three flavor QCD ($m_{\text{up}} = m_{\text{down}} \neq m_{\text{strange}}$ or short $N_f = 2 + 1$ QCD) without taking into account QED effects. More details of this analysis already appeared in [2].

The report is organized in the following way: in the next section we will present the data used in this analysis as well as the methods to obtain the decay-constant ratio f_K/f_π in the isospin symmetric limit without electromagnetic corrections. In Sec. 3 we then will convert our result to one for f_{K^\pm}/f_{π^\pm} and discuss the implications for the CKM-matrix elements V_{ud} and V_{us} .

2. Decay-constant ratio at the physical point

In this section we will first introduce the ensembles used for our analysis of the leptonic decay-constant ratio. Next we will discuss the functional forms used for the necessary inter- and extrapolations to physical quark masses and the continuum as well as the infinite-volume limits, respectively. Also the fit ranges applied to our data will be discussed before, eventually, our result for f_K/f_π is presented with all sources of uncertainties considered.

2.1 Ensembles used in the analysis

In total, 47 ensembles of gauge configurations generated by the Budapest-Marseille-Wuppertal collaboration (BMW-c.) with $N_f = 2 + 1$ flavors of tree-level clover-improved Wilson fermions and the tree-level Symanzik-improved gauge action were used. See [3, 4] for more details on the generation of these ensembles and other analyses carried out using them. The ensembles were generated at five different values for the gauge coupling β , resulting in lattice scales $1/a$ in the range from 1.7 to 3.7 GeV (using the Ω -baryon mass to set the scale). The (degenerate) light quark masses m_{ud} were such that pion masses between 130 and 680 MeV were obtained. In all but two cases the strange-quark mass was chosen close to its physical value. In Fig. 1 we show the kaon mass M_K as a function the pion mass M_π (left panel) as well as the product $(M_\pi L)$ as a function of M_π for the 47 ensembles used in our analysis. Here L is the spatial extent of the simulated lattice. More details and tables containing simulation parameters and measured quantities can be found in [2]. For details on how the decay constants of the pion and the kaon, f_π and f_K , resp., were obtained, we refer to [5].

¹ Here we adopt the convention, that the CKM-matrix elements V_{ud} and V_{us} are real and the complex phase in the first row only shows up in V_{ub} .

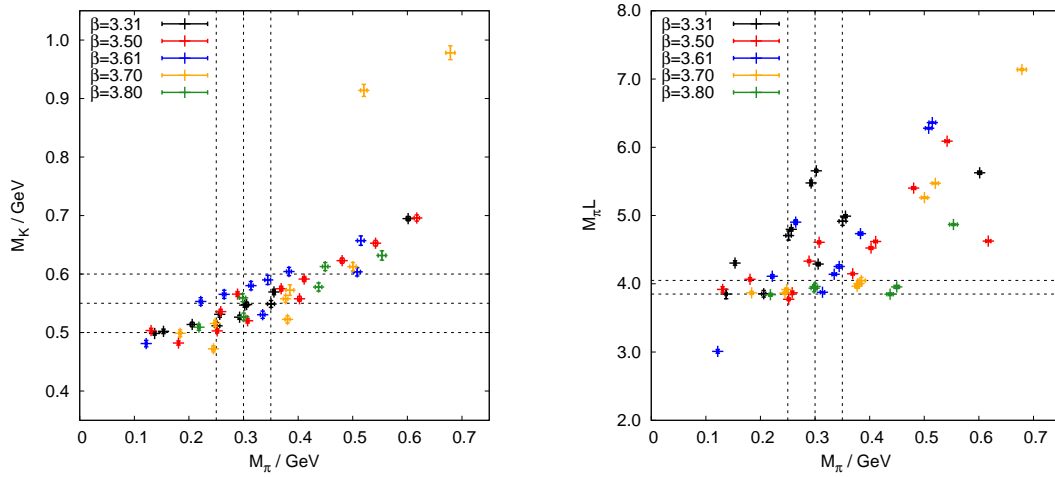


Figure 1: Overview of the ensembles used in this analysis. The *left panel* shows the kaon mass M_K as function of the pion mass M_π . To convert from lattice units to GeV, the mass of the Ω -baryon has been used to determine the lattice scale $1/a$ for each set of ensembles at the same gauge coupling β (mass-independent scale-setting). The *right panel* shows $(M_\pi L)$ as function of M_π . The *vertical and horizontal dashed lines* indicate the fit ranges used as discussed in Sec. 2.2

2.2 Functional fit forms and fit ranges

We have to interpolate our data to the physical quark masses as well as to extrapolate to the continuum and infinite-volume limits. In order to be able to quantify the systematic uncertainty associated with each of these inter- and extrapolations, we will always consider at least two different choices for the functional forms used and apply various fit ranges. Finally, we will take the variation of the results from the various fits as our systematic uncertainty.

For the interpolation to the physical mass point, given by $M_\pi^{\text{phys}} = 134.8(0.3) \text{ MeV}$ and $M_K^{\text{phys}} = 494.2(0.4) \text{ MeV}$ following the recommendation of the FLAG report [6]² which corrects the experimentally measured values M_{π^\pm} and M_{K^\pm} for strong isospin breaking and electromagnetic effects, either we use a fit form based on SU(3) chiral perturbation theory (ChPT) [8]

$$\frac{f_K}{f_\pi} = 1 + \frac{c_0}{2} \left\{ \frac{5}{4} M_\pi^2 \log \left(\frac{M_\pi^2}{\mu^2} \right) - \frac{1}{2} M_K^2 \log \left(\frac{M_K^2}{\mu^2} \right) - \frac{3}{4} M_\eta^2 \log \left(\frac{M_\eta^2}{\mu^2} \right) + c_1 [M_K^2 - M_\pi^2] \right\} \quad (2.1)$$

with two fit parameters c_0, c_1 and $M_\eta^2 = \frac{1}{3}(4M_K^2 - M_\pi^2)$ and μ being an arbitrary scale for the renormalization of ChPT. Alternatively, one of the following polynomial forms with either three, four, or six fit parameters $c_i^{n\text{-par.}}$ are used.

$$\frac{f_K}{f_\pi} = 1 + [M_K^2 - M_\pi^2] \left(c_0^{3\text{-par}} + c_1^{3\text{-par}} [M_K^2 - M_\pi^2] + c_2^{3\text{-par}} M_\pi^2 \right) \quad (2.2)$$

$$\frac{f_K}{f_\pi} = 1 + [M_K^2 - M_\pi^2] \left(c_0^{4\text{-par}} + c_1^{4\text{-par}} [M_K^2 - M_\pi^2] + c_2^{4\text{-par}} M_\pi^2 + c_3^{4\text{-par}} M_\pi^4 \right) \quad (2.3)$$

²At the time of writing the original article [2] the latest FLAG report [7] was not yet available to us. But except for a marginally change in $M_K^{\text{phys}} = 494.2(0.3) \text{ MeV}$ the values used in this work for M_π^{phys} and M_K^{phys} are not affected.

$$\begin{aligned} \frac{f_K}{f_\pi} = 1 + [M_K^2 - M_\pi^2] & \left(c_0^{6\text{-par}} + c_1^{6\text{-par}} [M_K^2 - M_\pi^2] + c_2^{6\text{-par}} M_\pi^2 + c_3^{6\text{-par}} M_\pi^4 \right. \\ & \left. + c_4^{6\text{-par}} M_\pi^2 [M_K^2 - M_\pi^2] + c_5^{6\text{-par}} [M_K^2 - M_\pi^2]^2 \right) \end{aligned} \quad (2.4)$$

Note that all these fit forms obey the flavor-symmetry constraint $f_K/f_\pi|_{m_{\text{ud}}=m_s} = 1$.

For the continuum limit extrapolation, we either assume scaling by a^2 or αa , where α is the strong coupling parameter at the simulated gauge coupling β . In each ansatz, one fit parameter appears, see [2] for the explicit expressions we applied. In order to also take into account the systematic uncertainty from setting the lattice scale, we adopted two different ansätze using the mass of the Ω -baryon (taken from [9]): a mass-independent one and an ansatz, where the scale was set per-ensemble, again for details we refer to [2] where also the values used for α are tabulated. The two extrapolation ansätze for the infinite-volume limit are based on the ChPT-formulae for the volume dependence [10, 11, 12, 13], where in each ansatz again one fit parameters appears, see [2]. These different choices already result in 32 combinations of fit forms with four, five, six, or eight fit parameters.

The fit ranges in the masses, lattice volumes, and lattice couplings were chosen as follows:

$$\begin{aligned} M_\pi &\leq 250 \text{ MeV}, 300 \text{ MeV}, 350 \text{ MeV}, \text{ no bound}, \\ M_K &\leq 500 \text{ MeV}, 550 \text{ MeV}, 600 \text{ MeV}, \text{ no bound}, \\ (M_\pi L) &\geq \text{no bound}, 3.85, 4.05, \\ \beta &\geq 3.31, 3.50, 3.61. \end{aligned}$$

These are also indicated by the horizontal and vertical dashed lines in Fig. 1. Together with the 32 considered combinations of fit forms and the requirement to only consider “true fits”, i.e. only fits with at least one degree of freedom, our data allowed for a total of 1368 different fits. For examples of specific fits, again, we refer to [2]. The combined results of these fits will be discussed in the following section.

2.3 Results

In Fig. 2 we show the p -values obtained in different fits, which were discussed in Sec. 2.2, plotted against the decay-constant ratio at the physical point (physical masses, infinite-volume limit, continuum limit) as obtained from a specific fit. In order to quantify the systematic uncertainty from our extrapolation to the physical point, we look both at the p -value weighted average and its variance and at an unweighted (“flat”) average, finding

$$\left. \frac{f_K}{f_\pi} \right|_{p\text{-value}} = 1.173(11)_{\text{stat}}(29)_{\text{syst}}, \quad \left. \frac{f_K}{f_\pi} \right|_{\text{flat}} = 1.191(08)_{\text{stat}}(24)_{\text{syst}}. \quad (2.5)$$

In order to try to reduce any bias introduced by our choice of fitting procedures etc. for our final value of the decay-constant ratio we take the straight average of these two resulting in

$$\frac{f_K}{f_\pi} = 1.182(10)_{\text{stat}}(26)_{\text{syst}} = 1.182(28)_{\text{comb}}. \quad (2.6)$$

The implications of this result for the CKM-matrix elements will be discussed in Sec. 3.

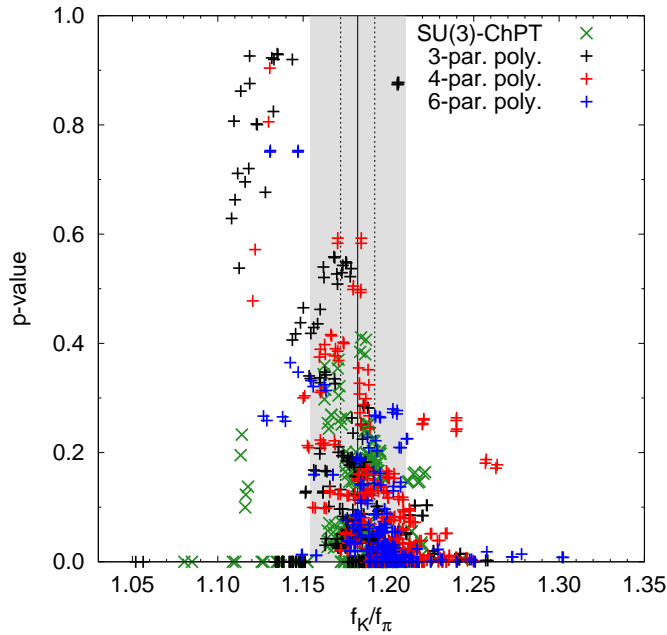


Figure 2: Scatter plot of the p -value obtained in each of the 1368 performed fits against the extrapolated (physical masses, infinite-volume and continuum limit) decay-constant ratio f_K/f_π from that particular fit. The *solid black line*, *dashed black lines*, and the *grey-shaded area* indicate the central value, statistical and combined error, resp., of our final result, Eq. (2.6).

We also looked at different subsets of our fits, sorted either by the fit form used for the mass interpolation, the continuum-limit extrapolation and so on or sorted by the fit ranges in M_π , M_K etc. The results are shown in the two panels of Fig. 3 (a detailed table of these results can be found in [2]). While most subsets in each group show comparable results, the biggest variation seems to stem from the variation in the cut to the lattice spacing a equivalent to the gauge coupling β (group at the bottom of the right panel of Fig. 3).

3. Discussion

Since in Eq. (1.1) the real-world ratio f_{K^\pm}/f_{π^\pm} appears, which also includes electromagnetic and $m_{\text{up}} \neq m_{\text{down}}$ effects, we have to correct our result for f_K/f_π , Eq. (2.6), to account for these effects. An analysis [8, 14] utilizing ChPT shows, that this correction can be obtained in the following way

$$\frac{f_{K^\pm}}{f_{\pi^\pm}} = \frac{f_K}{f_\pi} \sqrt{1 + \delta_{\text{SU}(2)}}, \quad (3.1)$$

where the correction term was estimated within ChPT [14] to have the value $\delta_{\text{SU}(2)} = -0.0043(12)$. The FLAG-report [6]³ quotes very similar values from lattice simulations with $N_f = 2 + 1$ flavors. Also, this correction was determined for $N_f = 2$ flavors in [15] to be $\delta_{\text{SU}(2)} = -0.0078(7)$. Conservatively, we decided to use the average of these two values with a 100 per cent uncertainty: $\delta_{\text{SU}(2)} = -0.0061(61)$ in Eq. (3.1) to correct our result, Eq. (2.6). We obtain

$$\frac{f_{K^\pm}}{f_{\pi^\pm}} = 1.178(10)_{\text{stat}}(26)_{\text{syst}} = 1.178(28)_{\text{comb}}. \quad (3.2)$$

³The same holds true for the values reported in the latest FLAG report [7] which appeared recently as a preprint.

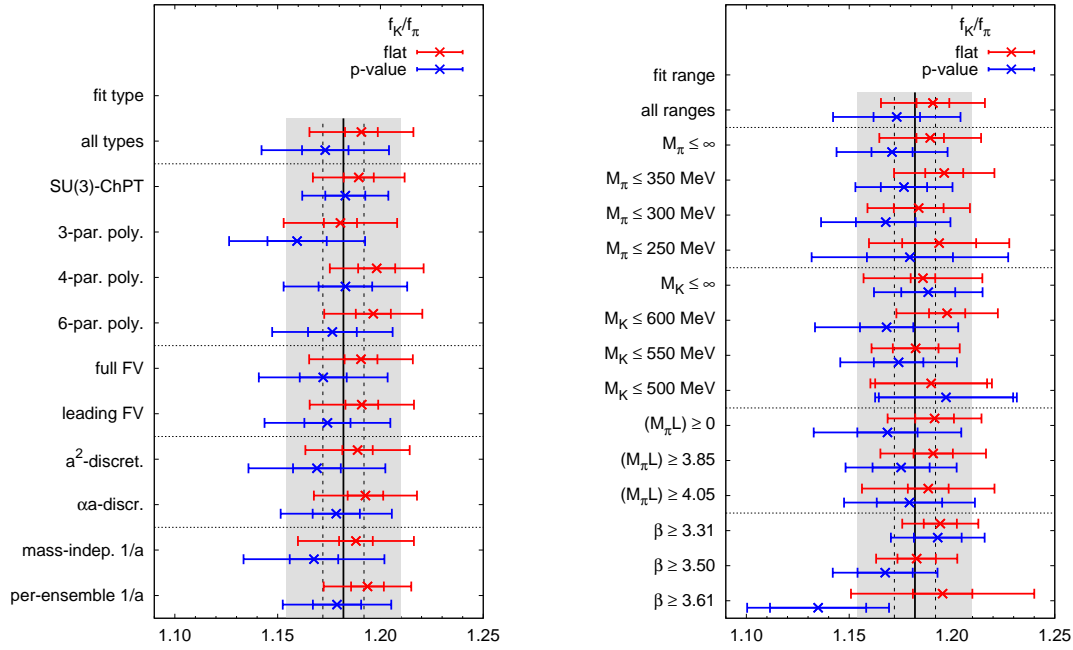


Figure 3: Results obtained without weighting (red symbols) and with p -value weighting (blue symbols) from various subsets of our fits. The *left panel* shows subsets given by the choices of fit forms, while the *right panel* shows subsets given by different fit ranges. The *solid black line*, *dashed black lines*, and the *grey-shaded area* indicate the central value, statistical and combined error, resp., of our final result, Eq. (2.6).

Now, we are able to use our value, Eq. (3.2), together with the experimental result [16, 17]

$$\frac{V_{us}}{V_{ud}} \frac{f_{K^\pm}}{f_{\pi^\pm}} = 0.27599(29)(24) = 0.27599(38)_{\text{exp}} \quad (3.3)$$

to determine the ratio of CKM-matrix elements

$$\frac{V_{us}}{V_{ud}} = 0.2343(20)_{\text{stat}}(52)_{\text{syst}}(03)_{\text{exp}} = 0.2343(55)_{\text{comb}}. \quad (3.4)$$

If one assumes first-row unitarity $V_{ud}^2 + V_{us}^2 + |V_{ub}|^2 = 1$ of the CKM-matrix together with the result from [17] that $|V_{ub}| = 4.12(37)(06) \cdot 10^{-3}$ is very small compared to the other two matrix elements, we obtain from our result, Eq. (3.4),

$$V_{ud} = 0.9736(04)_{\text{stat}}(11)_{\text{syst}}(01)_{\text{exp}}, \quad V_{us} = 0.2281(18)_{\text{stat}}(48)_{\text{syst}}(03)_{\text{exp}}. \quad (3.5)$$

Alternatively, one may use the high-precision result from super-allowed nuclear β -decays [18] $V_{ud} = 0.97417(21)_{\text{nuc}}$ to obtain from Eq. (3.4)

$$V_{us} = 0.2282(19)_{\text{stat}}(51)_{\text{syst}}(03)_{\text{exp+nuc}}. \quad (3.6)$$

and use these two results together with the above $|V_{ub}|$ to check for first-row unitarity

$$V_{ud}^2 + V_{us}^2 + |V_{ub}|^2 - 1 = 0.0011(09)_{\text{stat}}(23)_{\text{syst}}(05)_{\text{exp+nuc}} = 0.0011(25)_{\text{comb}}, \quad (3.7)$$

showing that within the errors the assumption of first-row unitarity of the CKM-matrix is fulfilled.

Comparing our result, Eq. (3.2), with the average from the latest FLAG-report [7] for $N_f = 2 + 1$ lattice-QCD simulations: $f_{K^\pm}/f_{\pi^\pm} = 1.192(5)$, one observes that the two agree within our large uncertainty, although our result favors a smaller decay-constant ratio. Definitely, our uncertainty needs to be reduced for one to be able to make a more robust statement. The same holds true if one compares our results for the CKM-matrix elements, Eq. (3.5), with the latest $N_f = 2 + 1$ FLAG-averages [7]: $V_{ud} = 0.97451(23)$, $V_{us} = 0.2243(10)$. For that reason, currently we are examining in which way different choices of fit ranges, fit forms might help to improve the precision of our result.

Acknowledgment We thank our colleagues Zoltán Fodor, Christian Hölbling, Stefan Krieg, Laurent Lellouch, Thomas Lippert, Thomas Rae, Andreas Schäfer, Kálmán K. Szabó, and Lukas Varnhost for the collaboration on this project.

Computations were carried out on the BG/Q supercomputer JUQUEEN at Forschungszentrum Jülich through a NIC grant, on Turing at IDRIS, France, under GENCI-IDRIS grant 52275, and on a local cluster at the University of Wuppertal. This work was supported by the DFG through the SFB/TRR 55 “Hadron Physics from Lattice QCD”, by the EU Framework Programme 7 grant (FP7/2007-2013)/ERC No 208740, by the OTKA grant OTKA-NF-104034, and by the projects OCEVU Labex (ANR-11-LABX-0060) and A*MIDEX (ANR-11-IDEX-0001-02). In addition, E.E.S. acknowledges support from the EU grant PIRG07-GA-2010-268367.

References

- [1] W. J. Marciano, Phys. Rev. Lett. **93**, 231803 (2004), arXiv:hep-ph/0402299.
- [2] S. Dürr *et al.*, (2016), arXiv:1601.05998.
- [3] S. Dürr *et al.*, Phys. Lett. **B701**, 265 (2011), arXiv:1011.2403.
- [4] S. Dürr *et al.*, JHEP **08**, 148 (2011), arXiv:1011.2711.
- [5] S. Dürr *et al.*, Phys. Rev. **D90**, 114504 (2014), arXiv:1310.3626.
- [6] S. Aoki *et al.*, Eur. Phys. J. **C74**, 2890 (2014), arXiv:1310.8555.
- [7] S. Aoki *et al.*, (2016), arXiv:1607.00299.
- [8] J. Gasser and H. Leutwyler, Nucl. Phys. **B250**, 465 (1985).
- [9] Particle Data Group, K. A. Olive *et al.*, Chin. Phys. **C38**, 090001 (2014).
- [10] M. Lüscher, Commun. Math. Phys. **104**, 177 (1986).
- [11] J. Gasser and H. Leutwyler, Phys. Lett. **B184**, 83 (1987).
- [12] G. Colangelo and S. Dürr, Eur. Phys. J. **C33**, 543 (2004), arXiv:hep-lat/0311023.
- [13] G. Colangelo, S. Dürr, and C. Haefeli, Nucl. Phys. **B721**, 136 (2005), arXiv:hep-lat/0503014.
- [14] V. Cirigliano and H. Neufeld, Phys. Lett. **B700**, 7 (2011), arXiv:1102.0563.
- [15] G. M. de Divitiis *et al.*, JHEP **04**, 124 (2012), arXiv:1110.6294.
- [16] M. Moulson, (2014), arXiv:1411.5252.
- [17] J. L. Rosner, S. Stone, and R. S. Van de Water, (2015), arXiv:1509.02220.
- [18] J. C. Hardy and I. S. Towner, Phys. Rev. **C91**, 025501 (2015), arXiv:1411.5987.